Network-based membrane filters: Influence of network and pore size variability on filtration performance☆

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ARTICLE INFO

ABSTRACT

We model porous membrane filters as networks of connected cylindrical pores via a random network generation protocol, and their initial pore radii via a uniform distribution of widths that vary about some mean value. We investigate the influence of network and pore size (radius) variations on the performance of membrane filters that undergo adsorptive fouling. We find that membrane porosity variations, independently of whether induced by variations of the pore radii or of the random pore network, are an important factor determining membrane filter performance. Network and pore size variations still play a role, in particular if pore radii variations are significant. To quantify the influence of these variations, we compare the performance metrics of networks built from pores of variable radii to their (equal porosity) counterparts built from pores of uniform radius. We show that the effect of pore radii variations is to increase throughput, but also to reduce foulant control.

1. Introduction

Membrane filtration is an important separation process used in many industrial and commercial applications such as treatment of radioactive sludge, water purification, beer clarification [1], semiconductor and microelectronics processing [2], air filtration [3] and membrane bioreactors [4]. Membrane filters used in these applications have a wide range of architectures, ranging from single layer thin porous films to multilayered porous membranes [5,6] to large scale continuous sheets of layered fibrous material [7–9].

Many models to describe the underlying membrane pore structure and/or geometry have been proposed and studied in recent years. For example, there are simple theoretical models to analyze the performance of membranes composed of multiple layers of different porous materials [6,10], or membranes with simple branched structures that can incorporate porosity gradients [11,12], among many others). As recent advances in imaging techniques have greatly contributed to the ability to compare such structural models to experiments ([13–16]; also see [17] for detailed images), more sophisticated models of membrane architecture have also been formulated, with a recent focus on accurate modeling of membrane filters with a network-type structure, for example, membranes where the solid component is comprised of fibers (so-called node–fibril type membranes) or those that transport

feed through networks of capillaries [18–20] (also used in the context of condensate banking [21].

The present paper focuses on the latter type of pore network model that captures the sponge-like structure of the membrane selective layer [22–24]. Our analysis of pore size variations builds on a model introduced previously in [20]. Briefly, a network model involves vertices and edges that represent pore junctions and throats respectively (referred to simply as *junctions* and *pores* from hereon respectively). Each pore is assumed to be a circular cylinder of fixed radius, with Hagen–Poiseuille flow, and conservation of fluid flux is imposed at each junction. Foulant is advected through pores by fluid and deposits on pore walls via adsorptive fouling. Membrane filter performance is analyzed by recording total throughput (volume of filtrate collected during the lifetime of the filter) and accumulated foulant concentration at the membrane outlet.

The primary findings of [20] that motivate the present work are the following relationships. Firstly, the initial porosity (pore void volume divided by total domain volume) of the network is demonstrated to be an important material feature that predicts total throughput by a power law, which holds particularly well when initial porosity > 0.5. Secondly, it was found that the accumulated foulant concentration at the membrane outlet decays exponentially with the initial tortuosity of the network (defined as the average distance traveled by a fluid

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particle from membrane inlet to outlet, see Appendix A.4 for a detailed definition).

In addition to the features discussed above, pore size (radius) variation is another important aspect of membrane filter design, which has been investigated by a number of authors, in particular regarding how it affects membrane selectivity and particle retention in a variety of applications [25–29]. It has not, however, been extensively studied in the context of network models for membrane filters: though such models allow for cylindrical pores with a distribution of lengths [19, 20], random variations of the pore radii have not, to our knowledge, been considered. Accordingly, in this paper, we focus on membrane filters whose pores may be considered as a network of interconnected capillaries of different initial radii. Henceforth in this work, the phrase "pore size variations" refers to variations in the initial pore radii, unless otherwise specified.

We present a novel characterization of the membrane pore network and model pore size variations via a set of random initial conditions for the radii of pores. Our assumption on the pore size distribution, in contrast to the log-normal distribution usually assumed in applications [30], has the advantage of avoiding extreme values and offering easier control over the variance, while still providing a reasonable description of effects such as inhomogeneities during the manufacturing process. Other pore size distributions (including the log-normal) tend to have input parameters that affect both the mean and variance simultaneously; furthermore, distributions that can potentially yield very large pore radius values would require more complicated fluid flow models than considered in this work.

We use our network representation to investigate the effect of pore size variations on the performance of membrane pore networks, as characterized by foulant concentration and total throughput of filtrate over the filter lifetime. We expect our results to improve our understanding of particle retention inside membrane pore networks and of how the rate of membrane fouling depends on input parameters. We note that other authors have taken complementary approaches to such issues; for example, see [31] for numerical comparisons of a number of pore size distributions, [32] for a study of the influence of pore size distribution on rheology in two-phase porous media flow, [33] for experimental discussions, and [34] for a discussion of related considerations in tissue engineering.

The paper is structured as follows: in Section 2, we describe the network model setup and introduce the key performance metrics. In Section 3, we set out our investigation strategy for pore size variation in an algorithm and declare the main nomenclature used in the analysis. In Section 4, we present and discuss our main results and in Section 5, we conclude our findings.

2. Setup: General pore networks

In this section we construct our model of a membrane filter represented by a random network of connected cylindrical pores. In Section 2.1, we describe how we generate a network that represents the internal pore structure of a membrane filter; in Sections 2.2 and 2.3, we outline the governing equations for the Hagen–Poiseuille fluid flow, for the advection of foulant particles carried by the flow, and for the pore evolution in time; in Section 2.4 we specify our notation; in Section 2.5, we provide the relevant physical scales; and in Section 2.6, we introduce two metrics to characterize the performance of a membrane filter with a specified pore network.

2.1. Network generation

We first demonstrate the 3D random network generation protocol employed in this work (following [20]). In Fig. 1, for illustration purposes, we show a 2D schematic of this protocol: the domain is a rectangular cell (square prism in 3D) of lateral height two times its horizontal width W, assumed to repeat periodically at the lateral
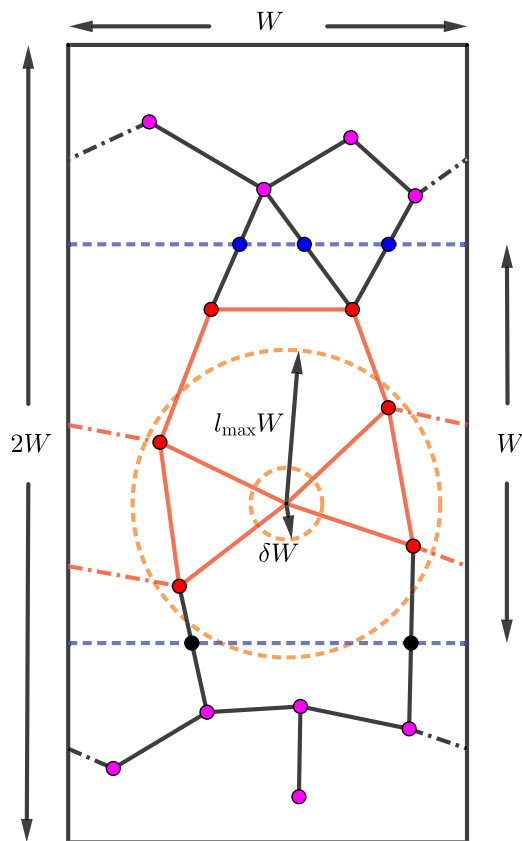

Fig. 1. 2D schematic of the 3D network generation with periodic boundary conditions showing: interior junctions $\mathcal{V}_{\text{int}}$ (red filled circles); pore inlets $\mathcal{V}_{\text{in}}$ (blue filled circles) and outlets $\mathcal{V}_{\text{out}}$ (black circles) induced by the cutting process; the cutting planes blue dashed lines; discarded points (magenta filled circles). Solid lines represent pores, while dash-dotted lines are pores that arise from the periodic boundary condition (red are interior to the membrane). $l_{\max}W$ and δW are prescribed maximum and minimum pore lengths respectively. (For interpretation of the references to color in this figure legend, the reader is referred to the web version of this article.)

boundaries (more details in Appendix A.1). We generate a pore network by first uniformly distributing N_{total} (a large integer to be prescribed) points as the *junctions* (nodes of the network) inside the domain. We construct *pores* (the *edges* of the network) by connecting all pairs of points that lie within a prescribed maximal distance $l_{\max}W$, but at least some minimal distance δW apart. We then cut the prism horizontally at two locations (see blue dotted lines in Fig. 1), and discard the end pieces to produce a square (or cube in 3D), which will represent an element of a membrane filter. The intersections of the two cutting planes and the pores (edges) naturally form membrane inlets and outlets, respectively (see Fig. 2 for a 3D realization, with lengths scaled by W).

2.2. Fluid flow

We next characterize fluid flow in the membrane network. The flow through the pores (which, within the current model, are cylinders of circular cross section) is assumed to obey the Hagen–Poiseuille model, valid provided the pores have sufficiently small aspect ratio (ensured in practice by choice of δ). The Hagen–Poiseuille equation states that fluid flux Q is proportional to pressure difference ΔP along each pore,

$$Q = K\Delta P, \quad K = \frac{\pi R^4}{8\mu L}, \tag{1}$$

where the conductance K depends on R and L, the radius and length respectively of the pore, and μ is fluid viscosity. To drive the flow, we
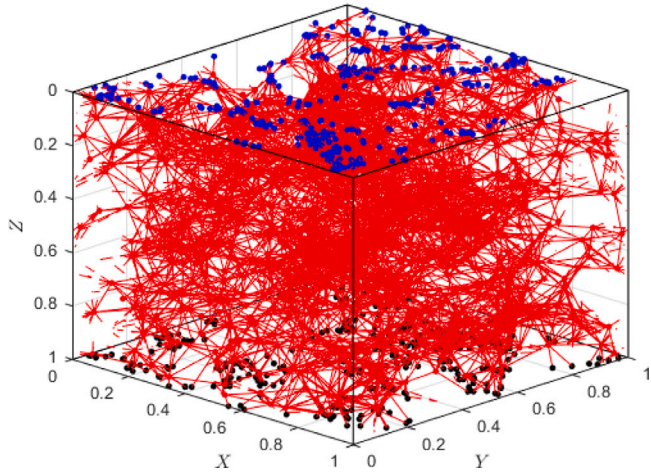
Fig. 2. Schematic of a 3D network with $l_{\max} = 0.15$ and $N_{\text{total}} = 2000$. Solid red lines are interior pores; dashed red lines are pores created by the periodic boundary conditions. Blue dots are inlets. Black dots are outlets.

prescribe the transmembrane pressure by setting the values on the top (P_0) and bottom (0) membrane surfaces, while the pressures at interior junctions are unknown. At each junction, we impose conservation of fluid flux, which leads to a system of equations (a discrete Laplace equation for these junction pressures, see Appendix A.2 for details). Once the pressures are found, flux through each throat is determined by Eq. (1).

2.3. Advection and adsorptive fouling

Feed solution enters the membrane top surface with a fixed foulant concentration (simply referred to as *concentration* from hereon), which serves as a boundary condition. Foulant particles are advected by the fluid while depositing on the pore walls, causing pores to shrink and eventually close up, a process known as adsorptive fouling. For each pore, let $C(Y, T)$ be the particle concentration at any point Y of the pore at time T, where Y is a local coordinate measuring distance along a pore in the direction of fluid flow. The concentration is assumed to satisfy the steady state advection equation (see [20] and references therein)

$$Q\frac{\partial C}{\partial Y} = -\Lambda R C, \quad 0 \le Y \le L, \tag{2}$$

$$C(0, T) := C_{\text{up}}(T). \tag{3}$$

Here, the right hand side of Eq. (2) is a sink of concentration that models foulant particles depositing on pore walls as they traverse the edge, with Λ (units of velocity) a parameter that captures the affinity between foulant particles and the membrane material. In reality Λ could change during filtration due to different affinities between the particles and the clean membrane, and the particles and the fouled membrane, the latter being more akin to particle–particle interactions; for simplicity, however, we assume Λ is a constant. In Eq. (3) $C_{\text{up}}(T)$ is the concentration at the upstream entrance of a pore. For pores connected to the membrane top surface (at inlets), we prescribe a constant concentration $C_{\text{up}}(T) = C_{\text{top}}$ as the upstream boundary condition, representing the concentration in the feed solution. The foulant concentrations at interior junctions are solved for using conservation of foulant particle flux QC; that is, the total combined incoming particle flux from upstream pores must be equal to that entering downstream pores. Similar to the fluid flux conservation, this law can be expressed by a system of equations that accounts for the network connectivity (see Appendix A.2 and [20]).

Table 1
Key dimensional quantities.

Dimensional quantity at time T	Symbol
Length of pore ij	L_{ij}
Radius of pore ij	$R_{ij}(T)$
Membrane unit length	W
Maximum pore length	$l_{\max}W$
Minimum pore length	δW
Pressure at junction i	$P_i(T)$
Concentration at junction i	$C_i(T)$
Flux in pore ij	$Q_{ij}(T)$
Deposition Coefficient	Λ

Lastly, as foulant particles deposit on the pore wall, pore radius decreases at a rate depending on local upstream concentration $C_{\text{up}}(T)$,

$$\frac{dR}{dT} = -\Lambda \alpha C_{\text{up}}(T), \quad R(0) = R_0, \tag{4}$$

where α is a parameter related to particle volume (see the appendix of [20] for details). The filter lifetime T_{final} is reached when the outgoing flux at the membrane's downstream surface falls to zero.

Eq. (4) implicitly assumes that the radius of each pore is a function only of time T and is independent of the local coordinate Y during its evolution, which is determined only by the upstream particle concentration. A more accurate model would be $\partial R(Y, T)/\partial T = -\Lambda \alpha C(Y, T)$; however, the simplification represented by Eq. (4) yields only minor differences in model outcome, while providing significant computational benefit since it permits analytical solution of Eq. (2), leading to immense computational speed-up (more than 100 times faster than the alternative). See [20] for further discussion, including a sufficient condition for this simplification to be reasonably valid. In particular, we have verified that the more accurate model leads to only a modest increase (at most 10%) in total throughput. The difference in accumulated foulant concentration is smaller still, and differences in the models' predictions of both performance measures are particularly insignificant when $l_{\max}$ (the ratio of the maximum pore length to the membrane thickness) is small. The computational savings afforded by this simplification facilitate the large number of simulations carried out in Section 3, needed to obtain reliable statistics.

2.4. Network notations

In Sections 2.2 and 2.3, we introduced fluid flow and adsorptive fouling in each individual pore and described the conservation laws that govern the physical quantities of interest in a pore network. Since describing the dynamics of such quantities requires coupling of all junctions and pores, we introduce the following indexing scheme. We use a single index $(\cdot)_i$ to indicate junction dependence only, e.g. pressure P_i at junction i. We use a double index $(\cdot)_{ij}$ to indicate dependence on the pore connecting junctions i and j. We do not distinguish between ij and ji for scalar quantities such as conductance K_{ij}; however, we do for quantities with a notion of direction, e.g., Q_{ij} as fluid flux from pore i to pore j. We summarize all physical quantities in Table 1.

2.5. Scales

We nondimensionalize the fluid flow and foulant transport model described in the previous sections as follows:

$$\begin{aligned} L_{ij} &= W l_{ij}, \quad (R_{ij}, R_{ij,0}) = W(r_{ij}, r_{ij,0}), \\ P_i &= P_0 p_i, \quad (C_i, C_{ij}) = C_{\text{top}}(c_i, c_{ij}), \\ \Lambda &= \frac{\pi W P_0}{8\mu} \lambda, \quad T = \frac{W}{\Lambda \alpha C_{\text{top}}} t, \end{aligned} \tag{5}$$

where the upper case symbols are dimensional quantities (listed in Table 1). Parameter values are listed in Table 2. We comment briefly

Table 2
Key dimensionless parameters.

Parameter	Symbol	Values
Initial number of junctions	N_{total}	500,775
Maximum pore length	$l_{\max}$	0.3
Minimum pore length	δ	0.06
Unperturbed pore radius	r_0	0.01
Noise Amplitude	β	0.06, 0.25, 0.5, 0.7
Number of networks	N_{net}	1000
Number of noise realizations	N_{noise}	500
Deposition coefficient	λ	5×10^{-7}

regarding the parameter λ: its main influence is in determining the time scale on which fouling occurs in the problem, and how deep within the pore network fouling particles penetrate. The numerical value of λ is expected to be small; such values promote more efficient use of the filter by allowing particles to permeate deeper into the membrane interior (larger λ values lead to a majority of fouling occurring very close to the membrane top surface). Our choice of the specific value, while somewhat arbitrary, satisfies this requirement. Further work, beyond the scope of this paper, would be needed to explore how the exact value assigned to this quantity influences the statistics of membrane performance metrics. We have confirmed numerically that moderate variations in λ do not change the trends that we describe in this work.

2.6. Performance metrics

We evaluate the performance of a membrane network using the following two metrics: (1) total throughput (h) and (2) accumulated foulant concentration at membrane outlet (c). Total throughput is the total volume of filtrate collected at the membrane outlet over the lifetime of the membrane network filter (at time t_{final}, when the flux through the membrane falls to zero). Accumulated foulant concentration measures the aggregate concentration of foulant particles in the collected filtrate when the filter is exhausted. Precise definitions of these quantities are given in Appendix A.3.

3. Investigation methods

We study pore size variations by prescribing a random initial condition for each pore radius, $r_{ij,0}$. More precisely, we consider perturbations to uniform pores by imposing a multiplicative noise,

$$r_{ij,0} = r_0 \left(1 + \epsilon_{ij}\right), \tag{6}$$

where $\epsilon_{ij} \sim \text{Unif}\,(-\beta, \beta)$ is a uniform random variable with noise amplitude $0 < \beta < 1$, independent for each pore. For the rest of this work, when we use the terms *perturbation* or *noise*, we are referring to Eq. (6). This assumption on the distribution of pore size variations has the advantage that we can directly control the width of the distribution via the input parameter β, making the analysis and interpretation of the results simpler.

We will study selected simulation outputs, f, as model parameters vary, with a primary focus on the influence of noise amplitude, β. Specifically, in this work f will be one of the following:

- h, total throughput (throughput),
- c, accumulated foulant concentration at membrane outlet (concentration),
- ϕ, initial network porosity (porosity),
- τ, initial tortuosity (tortuosity),

where the abridged terms in parentheses are used freely henceforth (see detailed definitions of h, c and τ in Appendices A.3 and A.4). We note also that since the nondimensional model operates in the domain of a unit cube, initial void volume and membrane porosity are equal in value. We refer to *porosity* only hereon.

A principal aim of this work is to compare the influence of two independent sources of randomness, namely, the random network generation process (*network variability* henceforth), and the random initial condition for the pore radius that yields pore radius variations (*noise variability* henceforth), on statistics of membrane filter performance metrics. In other words, network variability stems from variations in node and edge locations in membrane networks, while noise variability is due to pore size variations. Here we describe the methodology of our study, before summarizing the approach as an algorithm with enumerated steps below.

First, we generate a large number, N_{net}, of random membrane networks (per Section 2.1), each with the same initial pore radius r_0. We perturb the pore networks via Eq. (6) in the following two distinct ways: To probe noise variability, we fix one particular "typical" (to be made precise in what follows) network from the N_{net} that were generated, perturb it independently N_{noise} times (for N_{noise} sufficiently large), solve the governing Eqs. (1), (2) and (4) (coupled over the entire network using continuity as described) and collect statistics of performance metrics from these N_{noise} realizations of noise. To probe network variability, we perturb each of the N_{net} networks independently just once, and collect performance statistics from the perturbed networks.

Perturbing the pore radii inevitably changes the porosity of the network. This is important, since initial network porosity was shown to influence strongly both of our membrane performance metrics in the unperturbed case [20]. To investigate the importance of such induced porosity changes, we devise the following strategy. First, note that the initial porosity of a network is given by

$$\phi = \frac{\pi}{2} \sum_{ij} r_{ij,0}^2 l_{ij} \stackrel{\text{Eq. (6)}}{=} \frac{\pi}{2} r_0^2 \sum_{ij} \left(1 + 2\epsilon_{ij} + \epsilon_{ij}^2\right) l_{ij}, \tag{7}$$

where the last expression implies that perturbed networks have larger porosities on average, since ϵ_{ij}^2 is nonnegative (the factor of $1/2$ is needed because the double sum counts every pore twice). Note that this increase of porosity for perturbed network is influenced by the choice (uniform) of pore size perturbation distribution; other distributions (not considered in this work) that have nonzero mean may influence the porosity differently. For each perturbed network, we obtain its porosity via Eq. (7). We then impose this porosity on the underlying *unperturbed* network by determining a new initial pore radius such that the unperturbed and perturbed networks have the same porosity. This new unperturbed network is referred to as a *porosity-corrected* network, from which we also collect performance statistics. With these preparations, we define a 'score' as the difference between the outputs of the perturbed and porosity-corrected networks, normalized by the outputs of the porosity-corrected ones.

Each of the N_{net} membrane networks has a fixed maximum pore length $l_{\max} = 0.3$, minimum pore length $\delta = 0.06$ and initial unperturbed pore radius $r_0 = 0.01$. We vary two parameters: (1) noise amplitude β; and (2) total number of junctions N_{total}. The ensuing study first fixes an *average* porosity by fixing N_{total}, while varying β. All geometric parameters used in the algorithm are listed in Table 2, with their values.

We summarize the above procedures in the following algorithm:

1. (*Random network generation*) Choose N_{total}. Generate N_{net} unperturbed networks. Compute the initial porosity ϕ_0 for each network.
2. (*Noise perturbation*) Choose noise amplitude β.

    (a) (*noise variability*) Fix a typical network and perturb it N_{noise} times independently. Compute the outputs of the perturbed networks, f_{noise}, referred to as *output under noise realizations*. The typical network is chosen such that its porosity is the closest to the average porosity of the ensemble, a posteriori.

    (b) (*network variability*) Perturb each unperturbed network once independently via Eq. (6). Compute the associated outputs f_{net}, referred to as *output under network realizations*.

Table 3
Key outputs. The quantities in parentheses are the mean and standard deviation of the corresponding output.

Output	Symbol
Porosity-corrected	f_{pc}
Noise realizations	f_{noise} $(\overline{f_{noise}}, \sigma_{f_{noise}})$
Network realizations	f_{net} $(\overline{f_{net}}, \sigma_{f_{net}})$
Noise score	$\widehat{f}_{noise}$ $(\overline{\widehat{f}_{noise}}, \sigma_{\widehat{f}_{noise}})$
Network score	$\widehat{f}_{net}$ $(\overline{\widehat{f}_{net}}, \sigma_{\widehat{f}_{net}})$

Fig. 3. Scatter plot of throughput versus porosity, under (a) noise realizations, (b) network realizations (with each network perturbed once). The black rectangle in (b) shows the horizontal and vertical range of (a). For both plots, $\beta = 0.06$.

3. (*Porosity correction*) For each perturbed network, consider its underlying *unperturbed* equal-porosity network (created by prescribing the appropriate uniform pore radius for its pores; see Appendix A.5 for a short derivation). This new network is called a *porosity-corrected* network, with output labeled f_{pc}.

4. (*Scores*) We construct two scores that characterize the noise and network variability when comparing networks of equal porosity. The two scores are constructed independently from each other, so there is no confusion between notations.

   (a) (*Noise score*) We compute the following score,

$$\widehat{f}_{noise} = \frac{f_{noise} - f_{pc}}{f_{pc}}, \tag{8}$$

where f_{noise} and f_{pc} are computed via Steps 2a and 3, respectively. All quantities in Eq. (8) are vectors of length N_{noise}. The subtraction and division are element-wise.

   (b) (*Network score*) Perform Step 2b to obtain the outputs under network realizations f_{net}. Compute the following score,

$$\widehat{f}_{net} = \frac{f_{net} - f_{pc}}{f_{pc}}, \tag{9}$$

where f_{pc} are computed via Step 3. All quantities in Eq. (9) are vectors of length N_{net}. The subtraction and division are element-wise.

In all discussions below, we refer to the quantities defined by Eq. (8) or Eq. (9) as *porosity-corrected* scores.

5. Obtain the means and standard deviations of model outputs, $\overline{f_{noise}}$ and $\overline{f_{net}}$, and scores, $\overline{\widehat{f}_{noise}}$ and $\overline{\widehat{f}_{net}}$. The means under noise and network realizations are computed by averaging over the number of noise (N_{noise}) and network (N_{net}) realizations, respectively.

6. Go back to Step 2 with a different β.

7. Go back to Step 1 with a different N_{total} (to vary initial porosity).

All outputs computed in this algorithm are summarized in Table 3.

4. Results and discussions

In Section 4.1, we study the performance metrics under a specific choice of model parameters as an example. First, we compare the noise and network variability of the raw metrics (throughput and concentration) and geometric quantities (porosity and tortuosity). Then, we compare the porosity-corrected scores under noise (per Eq. (8)) and network (per Eq. (9)) realizations (steps 4a and 4b in the algorithm, respectively). In Section 4.2, we reinforce the example by a thorough sweep of the parameter space and present our main results.

4.1. Detailed example: Low porosity network in low noise regime

In this section, we present a set of results for membrane pore networks in the regime of low noise amplitude perturbations to the pore radii and low initial porosity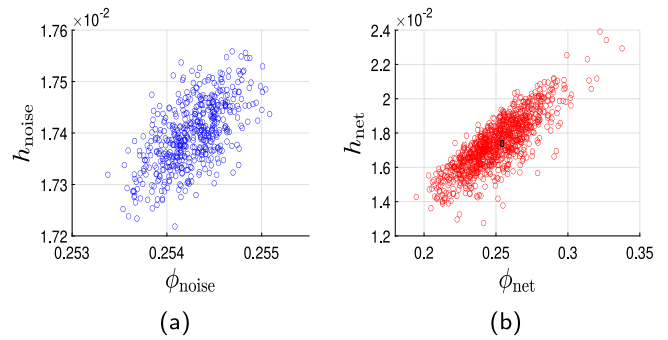. We choose this parameter regime as a detailed example because the results at higher porosities are qualitatively similar but more time-consuming to compute. We generate N_{net} networks with an initial number of points N_{total}, which yield an *ensemble average* initial porosity $\phi \approx 0.25$ (averaged over N_{net} unperturbed networks). The noise ϵ_{ij} is realized N_{noise} times for each network, with fixed noise amplitude $\beta = 0.06$ here. We give values of N_{net}, N_{noise} and N_{total} in Table 2 and have found that these numbers are sufficient to account for the random nature of the network and noise generation protocol; larger numbers produce similar results.

When we study noise variability (per Step 2a in the algorithm), we fix a typical network with initial porosity very close to the ensemble average, 0.25. We first compare the statistics of performance metrics (per Steps 4a and 4b in the algorithm) and then discuss the similarities and differences in the network and noise scores calculated for the initial porosity and tortuosity of the networks, similarly averaged over many network and noise realizations. In particular, we focus on the two relationships found in [20] as example results: throughput vs. porosity, and concentration vs. tortuosity.

4.1.1. Throughput and porosity

Fig. 3 shows throughput versus porosity under (a) noise and (b) network realizations, respectively. Both figures show that, as expected, throughput is an increasing function of porosity because more porous filters process more filtrate. However, we note the different scales in the two figures: in Fig. 3(b) we also plot a black rectangle that represents the total range of Fig. 3(a) showing that, before we correct for induced porosity changes, network variations incur much more variation in total throughput than pore size variations in this noise/porosity regime. In other words, the variability induced by noise is significantly smaller than that induced by network variability.

As noted in Section 3 (per Eqs. (6) and (7)), and as evident from Fig. 3, each pore size perturbation leads to a porosity variation. Since initial porosity is known to be a key parameter determining filter performance [20], we proceed by considering the *noise* and *network scores* that correct for induced porosity changes, formulated in Eqs. (8) and (9) in Step 4 of the algorithm in Section 3. Fig. 4 presents histograms of throughput scores under noise and network realizations. We see that the two histograms are very similar. Comparing Figs. 3 and 4, we deduce that the porosity change is the crucial factor (at least for the present (ϕ, β) values) since Fig. 4, in marked contrast to Fig. 3, shows that variations of network and noise as measured by the scores have a very similar effect.

4.1.2. Concentration and tortuosity

We next investigate the influence of noise perturbation on the accumulated concentration of particle impurities in the filtrate (concentration) and the tortuosity of the pore network (see Appendix A.4 for a detailed definition), quantities that are strongly related in unperturbed pore networks [20].
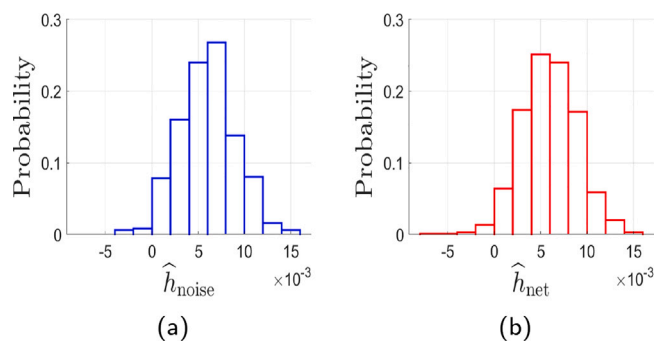
Fig. 4. Histogram of throughput score, under (a) noise realizations, (b) network realizations (with each network perturbed once). Same parameters as in Fig. 3.
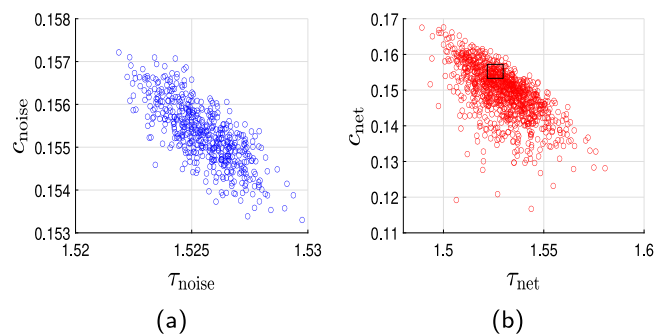

Fig. 5. Scatter plot of concentration versus tortuosity. Same description and parameters as in Fig. 3.
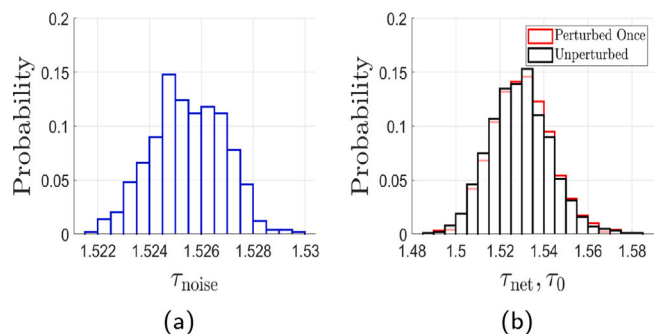

Fig. 6. Histogram of tortuosity under (a) noise realizations (b) network realizations. Same parameters as in Fig. 3; same data as Fig. 5.

Fig. 5 shows concentration versus tortuosity under (a) noise and (b) network realizations. First, both figures show that concentration is a decreasing function of tortuosity. This makes sense because the longer foulant particles travel in the network, the more likely they are to adsorb to the pore walls, leading to lower concentration at the membrane outlet. Second, we point out that the black rectangle in Fig. 5(b) covers the total range of Fig. 5(a), showing, once again, that before induced porosity changes are accounted for, network variability dominates over noise, for the present choice of parameters.

Figs. 6(a) and 6(b) plot the data of Fig. 5 as a histogram of tortuosity under noise and network realizations (respectively). Note the different horizontal ranges on the two plots (cf. Fig. 5). Also, in Fig. 6(b), we see that the distribution of tortuosity shifts slightly to the right relative to the unperturbed case, possibly due to some subtle geometric influence from the random network generation, not yet understood.

Fig. 7 shows histograms for the porosity-corrected concentration and tortuosity scores under noise (Figs. 7(a) and 7(c), generated using Eq. (8)) and network (Figs. 7(b) and 7(d), generated using Eq. (9))
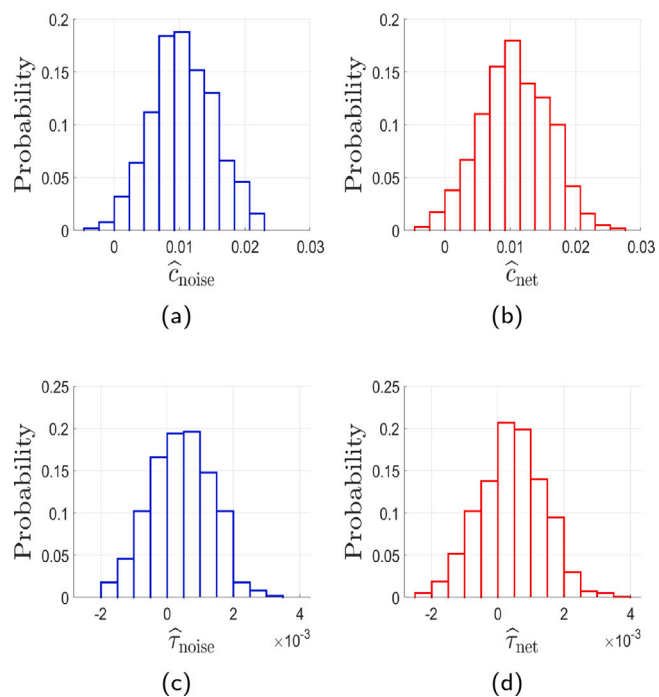

Fig. 7. Histogram of concentration and tortuosity scores, respectively under (a,c) noise realizations and (b,d) network realizations. Same parameters as in Fig. 3.

realizations. The histograms of concentration scores in Figs. 7(a) and 7(b) are very similar in shape and width, suggesting that, after we correct for porosity differences, perturbing one network many times is equivalent to perturbing many networks once.

4.2. Results for varied noise amplitude and porosity

We now explore the influence of noise amplitude, β, and of initial network porosity, ϕ. For brevity, we condense results for different initial porosity regimes in terms of the *average* outputs (see Step 5 of the algorithm) and their standard deviations (shown via error bars; see Table 3 for their notations) plotted as functions of β.

Figs. 8(a) and 8(b) plot the raw mean throughputs $\overline{h_{\text{noise}}}$ and $\overline{h_{\text{net}}}$ against noise amplitude β for low and high network porosity regimes respectively. These raw mean throughputs are increasing functions of β, an observation that we attribute largely to the fact that an increase of β on average increases initial porosity, which then allows more filtrate to be processed. Figs. 8(c) and 8(d) plot the mean throughput scores that correct for porosity changes, for filters of low and high porosity, respectively. We find that the mean throughput scores are also increasing functions of β; however, this increase (which is still significant, around 60% for $\beta = 0.75$ for both low and high porosity networks per Figs. 8(c) and 8(d)) is due purely to the pore size variations. This implies that, for equal porosity networks, when the magnitude of the noise perturbation β is large enough, pore radius variations do lead to an appreciable change in total throughput. This further suggests that pore radius variations promote filtrate production, at least for the considered uniform pore size distribution.

We now study the variability of throughput and throughput scores induced by pore radius variations, represented by the error bars (standard deviation of each statistic) for all cases shown in Fig. 8. To facilitate this comparison, we plot in Fig. 9 the ratio of the standard deviations of the mean quantities under noise and network realizations (referred to simply as the s.d. ratio below) from Fig. 8, for raw throughputs and their scores. A value of the s.d. ratio near 1 implies that the compared quantities have similar standard deviations, whereas a value
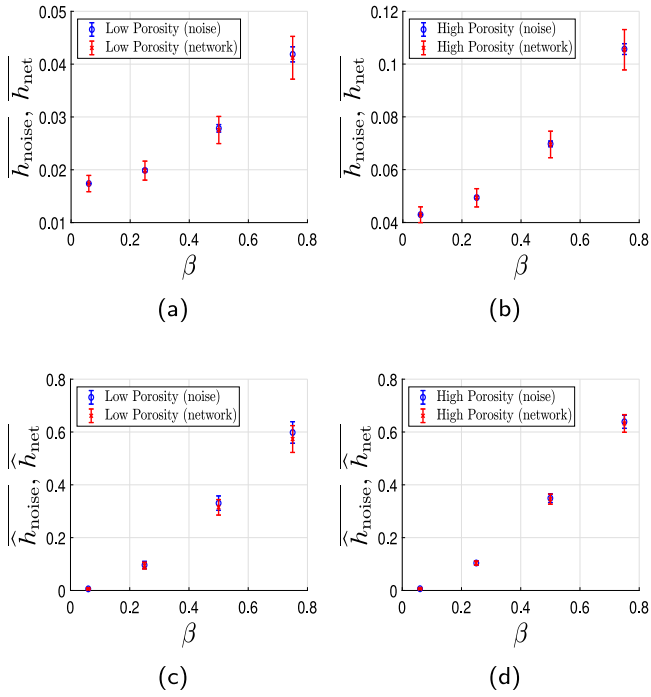
Fig. 8. Mean throughput $\overline{h_{\text{noise}}}$, $\overline{h_{\text{net}}}$ (a,b) and mean throughput scores $\widehat{\overline{h}}_{\text{noise}}$, $\widehat{\overline{h}}_{\text{net}}$ (c,d) versus noise amplitude β under noise and network realizations. Low porosity $\phi \approx 0.25$ (a,c) and high porosity $\phi \approx 0.6$ (b,d) cases are shown. Vertical error bars are standard deviations for each mean value. Note that the scores (c,d) show relative change, so that, e.g., an increase from 0 to 0.6 on the vertical axes means a 60% increase.
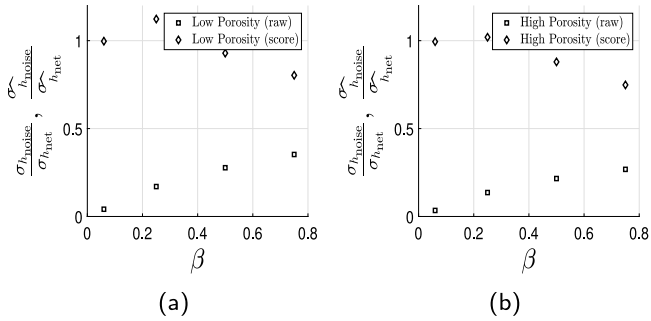

Fig. 10. Mean concentration and concentration score. Same setup as in Fig. 8.

Fig. 9. Ratio of standard deviations of raw throughput (black squares) and throughput scores (black diamonds) for (a) low porosity (corresp. Figs. 8(a) and 8(c)) and (b) high porosity (corresp. Figs. 8(b) and 8(d)).
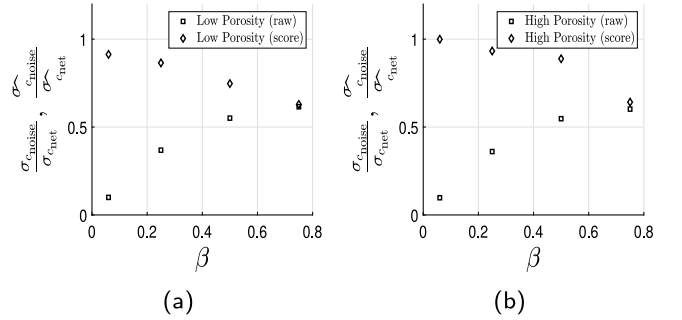

Fig. 11. Same setup as Fig. 9 for raw concentration and concentration scores. (a) low porosity (corresp. Figs. 10(a) and 10(c)) and (b) high porosity (corresp. Figs. 10(b) and 10(d)).

near 0 says that variation induced by network realizations dominates that from noise. In both low and high porosity regimes (Figs. 9(a) and 9(b) respectively), we observe that the s.d. ratios for the raw throughput (black squares) are all below 0.5 for any value of noise amplitude β, thus the throughput variability from network realizations dominates that from noise. However, the s.d. ratios of the throughput scores (black diamonds) are close to 1 in both porosity regimes, for all β values considered. This finding shows that, when considering the total throughput metric, network porosity accounts for the differences between network and noise variability.

Figs. 10(a) and 10(b) show mean concentrations $\overline{c_{\text{noise}}}$ and $\overline{c_{\text{net}}}$ of particles in the filtrate as a function of noise amplitude β for low and high porosity regimes. We observe that the mean raw concentrations in both cases are increasing functions of β. We attribute this increase largely to the increase of porosity with β, as evident in Eq. (7), allowing more particles to pass through unfiltered and thus worsening foulant control. In contrast, Figs. 10(c) and 10(d) plot the mean concentration scores $\widehat{\overline{c}}_{\text{noise}}$ and $\widehat{\overline{c}}_{\text{net}}$ (see Eqs. (8) and (9)), which correct for porosity

changes: the increase with β seen here is a direct consequence of the pore size variations. For example, for $\beta = 0.75$, the perturbed networks show more than 150% increase in concentration when compared with porosity-corrected unperturbed counterparts. This implies that pore radius variations influence foulant concentration in the filtrate significantly for sufficiently large noise amplitude β, and in fact are detrimental for foulant control.

Now, we discuss the observed variability of the raw mean concentration and concentration scores induced by pore size variations. Similar to the study of throughput variability, Fig. 11 plots the ratios of standard deviations of the mean concentrations (raw values and scores) under noise and network realizations, found in Fig. 10. The observed trends are similar to Figs. 9(a) and 9(b): before correcting for induced porosity changes, network variability dominates noise variability for all porosities and noise amplitudes tested, exemplified by the s.d. ratios for the raw concentrations (black squares). However, this dominance is weaker for particle concentration than for throughput: in Fig. 11 the s.d. ratios for the raw concentrations are larger than those for the raw throughputs in Fig. 9. This means that concentration, as a performance metric, experiences larger variations when pore radii are perturbed than does throughput. We further note that the s.d. ratio
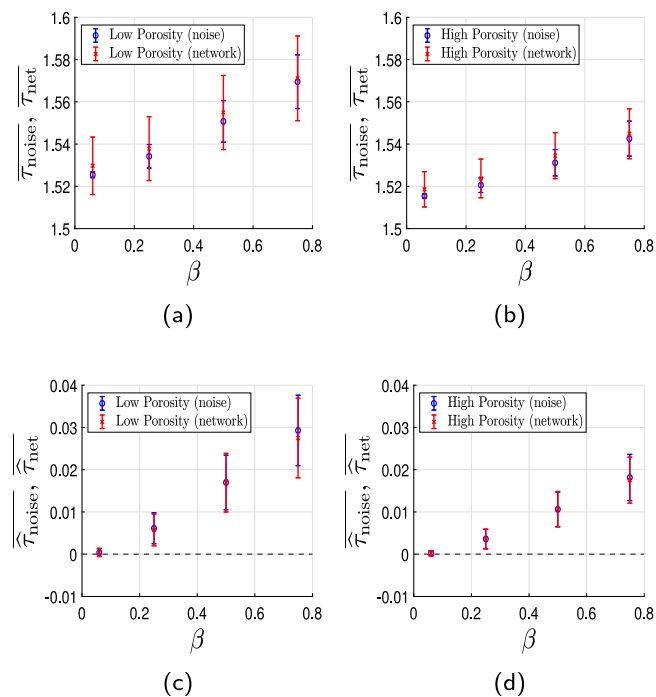
Fig. 12. Mean tortuosity and tortuosity score. Same setup as in Fig. 8.

of concentration scores (black diamonds) is close to 1 for small β (underscoring the importance of porosity), but decreases as β increases. In other words, porosity changes due to network variations are solely responsible for concentration variability at small noise amplitudes, but not at high amplitudes.

As noted earlier, tortuosity (the average dimensionless path length of a fluid element as it passes through the pore network; see definition in Appendix A.4) is an important quantity determining membrane filter performance [20]. Fig. 12 shows mean tortuosities $\overline{\tau_{\mathrm{noise}}}$ and $\overline{\tau_{\mathrm{net}}}$, and mean tortuosity scores $\widehat{\overline{\tau}}_{\mathrm{noise}}$ and $\widehat{\overline{\tau}}_{\mathrm{net}}$ as noise amplitude varies. We find that in both porosity regimes, network variability (red error bars) still dominates noise variability (blue error bars) if we do not correct for induced porosity variations. The dominance is less strong for higher noise amplitudes, however, where the error bar sizes become similar, indicating that sufficiently high noise amplitude induces similar variability in network tortuosity to that due to the network generation protocol. Figs. 12(c) and 12(d) show the means of the tortuosity scores $\widehat{\overline{\tau}}_{\mathrm{noise}}$ and $\widehat{\overline{\tau}}_{\mathrm{net}}$; we observe that after correcting for porosity changes, noise and network variability are now very similar for all noise amplitudes considered, in both porosity regimes. This implies that, when porosity variations are accounted for, the effects of noise and network variability on tortuosity are comparable.

Furthermore, we observe that the mean tortuosity and tortuosity scores in Fig. 12 are increasing functions of noise amplitude β in all cases, but the changes are modest compared to those in throughput seen in Fig. 8 (see a further discussion in Appendix A.4 on how β affects τ). Overall, tortuosity variations do not exceed 4% for all parameters considered (note the vertical scales of Figs. 12(c) and 12(d)). Based on this observation, we infer the relative influence of throughput and tortuosity scores on concentration scores. On the one hand, filters with larger throughput scores tend to allow more foulants to pass through and thus produce *higher* foulant concentrations in the filtrate. On the other hand, filters with larger tortuosity lead to *lower* concentrations [20]. Hence, as β increases, the increased mean throughput and tortuosity scores have opposite effects on mean concentration scores. The fact that we observe still an increased concentration score as β increases (per Figs. 10(c) and 10(d)) implies that the effect of increasing

throughput dominates that from increasing tortuosity. In other words, although pore size variations do increase pore network tortuosity (favorable for foulant control), this effect is apparently insufficient to lead to improved foulant removal when balanced against the tendency of the increased throughput to transport a greater proportion of foulant particles through the filter.

Before closing, we note in passing some perhaps unexpected features of the results. First, comparing Figs. 12(c) and 12(d), we observe that filters with higher porosity show a smaller increase in tortuosity with β. Second, we note that, according to Figs. 12(c) and 12(d) for the smallest value of β used, some networks may have *smaller* tortuosity than their unperturbed counterparts, as shown by the first error bar crossing the zero line.

5. Conclusions

In this work, we have devised and implemented a network model for adsorptive fouling of a membrane filter whose pore structure is formed by randomly generating junctions and pores. Within this network representation, we have modeled pore size variations as random initial conditions (uniform multiplicative noise for the pore radii). We have studied these two sources of randomness – the random network generation protocol and pore radius variations – by comparing their effects on two key membrane filter performance metrics: total throughput and accumulated concentration of adsorptive foulants in the filtrate.

We report three principal findings. First, we find that the initial porosity of the pore network is a critical feature of the filter, and a strong determinant of performance. Secondly, we observe that the influence of pore radius variations on membrane performance becomes prominent when the noise amplitude is large. Lastly, we conclude that pore size (radius) variations are favorable for maximizing filtrate production, but unfavorable for foulant control.

To arrive at these conclusions, we first compare membrane network performances directly by varying noise amplitude. Then, to rule out the effect from variations in porosity, we formulate the corresponding porosity-corrected networks and study the relative change in performance metrics via a well-defined score.

We note that the reported findings were obtained by assuming a uniform distribution of pore radius variations, and furthermore we have not considered large-particle sieving. Considering sieving in addition to adsorptive fouling and exploring the influence of other pore size distributions will be a focus of future work.

CRediT authorship contribution statement

B. Gu: Conceptualization, Methodology, Software, Validation, Formal analysis, Investigation, Resources, Data curation, Writing – original draft, Writing – review & editing, Visualization. **L. Kondic:** Conceptualization, Methodology, Validation, Formal analysis, Investigation, Resources, Data curation, Writing – review & editing, Supervision, Project administration, Funding acquisition. **L.J. Cummings:** Conceptualization, Methodology, Validation, Formal analysis, Investigation, Resources, Writing – review & editing, Supervision, Project administration, Funding acquisition.

Declaration of competing interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

Acknowledgments

Appendix

A.1. Network generation

We generate a membrane pore network via a variant of the Random Geometric network. To generate the set of pore junctions $\mathcal{V}$, we place N_{total} randomly distributed points (following a uniform distribution) in a rectangular box of height $2W$, with square cross-section of length W. We connect points that lie within a distance of $l_{\max} W$, but also at least δW apart, where δ is a control parameter (fixed throughout this work) such that when chosen large enough, it ensures validity of the Hagen–Poiseuille framework used to model fluid flow. These connections form a set of pores $\mathcal{E}$ and together with the junction set $\mathcal{V}$ we obtain an initial network $G = G(\mathcal{V}, \mathcal{E})$. We say $(v_i, v_j) \in \mathcal{E}$ when two junctions $v_i, v_j \in \mathcal{V}$ form a pore.

We then cut through the rectangular box with two horizontal planes at heights $0.5W$ and $1.5W$, respectively. The intersections of these two planes and the set of pores $\mathcal{E}$ form the set of inlets $\mathcal{V}_{\text{in}}$ and outlets $\mathcal{V}_{\text{out}}$ respectively. Altogether, the above procedure forms the domain for fluid flow and fouling, described in Section 2.

A.2. The graph Laplacian

We associate each network G with a (weighted) graph Laplacian, a generalization of the finite difference discretization of the classic Laplace operator $\nabla \cdot \nabla$. It is a square matrix whose off-diagonal terms indicate connection weights, and whose diagonal terms record the total weights of neighbors of each discretization point (junction). In our work, the most relevant weight is the conductance K_{ij} of each pore, given by

$$K_{ij} = \frac{\pi R_{ij}^4}{8\mu L_{ij}}.$$

Then the K-weighted graph Laplacian is defined as

$$\mathcal{L}_K := D - K, \tag{10}$$

where

$$D_{ij} = \begin{cases} \sum_{l=1}^{|\mathcal{V}|} K_{il}, & j = i, \\ 0, & \text{otherwise}, \end{cases} \tag{11}$$

where $|\mathcal{V}|$ is the number of junctions.

While the above setup characterizes the flux inside an individual pore, we employ conservation of flux at each interior vertex v_i throughout the network,

$$0 = \sum_{v_j : (v_i, v_j) \in E} Q_{ij}. \tag{12}$$

Combining Eqs. (1) and (12), we form a graph Laplace equation for the pressures P at each vertex, to which we add the specified pressure drop boundary conditions,

$$\mathcal{L}_K P(v_i) = 0, \quad v_i \in \mathcal{V}_{\text{int}}, \tag{13}$$

$$P(v) = P_0, \quad v \in \mathcal{V}_{\text{in}}, \tag{14}$$

$$P(v) = 0, \quad v \in \mathcal{V}_{\text{out}}. \tag{15}$$

Once the pressures $P(v_i)$ are found for all interior junctions $v_i \in \mathcal{V}_{\text{int}}$, we use Eq. (1) to find flux Q_{ij} in each pore to form a flux matrix $\mathbf{Q}$ with i and j as row and column indices respectively.

Using conservation of particle flux at each junction, we arrive at the following advection Laplace equation for foulant concentration $C_i(T)$ at each vertex $v_i \in \mathcal{V}_{\text{int}}$,

$$\mathcal{L}_{\mathbf{Q}}^{\text{in}} C = (\mathbf{Q} \circ B)^{\mathrm{T}} C_0, \quad B_{ij} = \exp\left(\frac{-\Lambda R_{ij} L_{ij}}{Q_{ij}}\right), \tag{16}$$

$$C_0 = (C_{\text{top}}, \ldots, C_{\text{top}}, 0, \ldots, 0)^{\mathrm{T}}, \tag{17}$$

where $\mathcal{L}_{\mathbf{Q}}^{\text{in}} = D_{\mathbf{Q}^{\mathrm{T}}} - (\mathbf{Q} \circ B)^{\mathrm{T}}$ is the *advection Laplacian* with a sink B, whose form arises from an analytical solution to Eq. (2). T and $\circ$ denote matrix transpose and the element-wise multiplication respectively. See [20] for a detailed derivation of this linear system.

A.3. Performance metrics

Volumetric throughput of a membrane filter over its operational lifetime is a commonly-used measure of overall efficiency. The volumetric throughput $H(T)$ through the filter is defined by

$$H(T) = \int_0^T Q_{\text{out}}(T')\, dT', \tag{18}$$

$$Q_{\text{out}}(T) = \sum_{v_j \in \mathcal{V}_{\text{out}}} \sum_{v_i : (v_i, v_j) \in E} Q_{ij}(T), \tag{19}$$

where $Q_{\text{out}}(T)$ is the total flux exiting the filter. With the scales chosen in Eq. (5), throughput H has scale

$$H = \frac{W^3}{\alpha C_0} h. \tag{20}$$

In particular, we compute $h_{\text{final}} := h(t_{\text{final}})$, the total volume of filtrate processed by the filter over its lifetime.

To connect with experiments and applications, we briefly discuss the order of magnitude of the dimensional throughput H in Eq. (20). The parameter α is of the order of foulant (contaminant) particle volume while C_0 is the number of particles per fluid (solvent) volume. The product αC_0 then yields an estimate of the volume ratio of the (contaminant) solute and the (fluid) solvent. For example, the permissible exposure limit by OSHA (Occupational Safety and Health Administration) for 1-dioxane in contaminated water is at a concentration of about 100 mg/L [35]. The density of 1-dioxane is close to that of water, and thus this level of concentration translates to a volume ratio (αC_0) of 10^{-4}. With this estimate and the order of magnitude for dimensionless throughput H seen in Fig. 3(a) (about 10^{-2}), H then is of the order of $10^2 W^3$ where W^3 represents the volume of a cubic unit of a membrane filter of side length W. This estimate suggests that this cubic unit will process filtrate volume of order 100 times its membrane material volume, given the parameters used in this work.

Another performance measure is the accumulated foulant concentration in the filtrate, which captures the aggregate particle capture efficiency of the filter. The accumulated foulant concentration is defined by

$$C_{\text{acm}}(T) = \frac{\int_0^T C_{\text{out}}(T')\, Q_{\text{out}}(T')\, dT'}{\int_0^T Q_{\text{out}}(T')\, dT'},$$

where

$$C_{\text{out}}(T) = \frac{\displaystyle\sum_{v_j \in \mathcal{V}_{\text{out}}} \sum_{v_i : (v_i, v_j) \in E} C_j(T)\, Q_{ij}(T)}{Q_{\text{out}}(T)}.$$

C_{acm} has scale $C_{\text{acm}} = C_{\text{top}} c_{\text{acm}}$. Of particular interest is $c_{\text{final}} := c_{\text{acm}}(t_{\text{final}})$, which provides a measure of the aggregate particle capture efficiency of the filter over its lifetime.

We further simplify the notations h_{final} and c_{final} to h and c in the main text as we consider only the end states of these performance measures in our analysis.

A.4. Tortuosity

Tortuosity τ of a membrane network is defined by the average distance traveled by a fluid particle from membrane top surface to bottom, relative to membrane thickness W. We here provide a formula via a probabilistic approach,

$$\tau = \frac{\pi_0^{\mathrm{T}}}{W} \left(\sum_{n=1}^{m} \mathbf{P}^{n-1} \right) \text{diag}\left(\mathbf{PW}_E \right), \tag{21}$$

where T means vector transpose and diag means the diagonal component of a matrix. Here we provide some intuition for each term. The initial distribution π_0 describes the probability of the fluid particle choosing an inlet on the membrane top surface. To calculate π_0, we compute the proportion of flux entering each inlet on the upstream surface relative to total flux. **P** within the sum describes the law of how a fluid particle traverses the network from one junction to its adjacent neighbors (known as *a step*); the upper limit m is the largest number of steps a particle takes to exit the membrane bottom surface, which can be found for each network. Lastly, $\mathrm{diag}\left(\mathbf{PW}_E\right)$ describes the average distance traveled by the fluid particle in one step starting from each junction. We refer the reader to [20] for details of the derivation.

In our study, since we perturb each pore radius via Eq. (6), tortuosity τ implicitly depends on the noise amplitude β through π_0 and **P** because they both involve the fluxes (affected by the pores' perturbed conductances).

A.5. Porosity correction

In Step 3 (*porosity correction*) of the algorithm, we derive the expression of r_{pc} such that $\phi_{\mathrm{noise}} = \phi_{\mathrm{pc}}$. It relies on writing ϕ_{pc} in terms of ϕ_0 (see List 1 and consider Eq. (7) with $\beta = 0$), the porosity of the unperturbed network with initial radius r_0:

$$
\begin{aligned}
\phi_{\mathrm{noise}} = \phi_{\mathrm{pc}} &= \frac{\pi}{2} r_{\mathrm{pc}}^2 \sum_{\mathrm{edge}} (\text{edge length}) \\
&= \left(\frac{\pi}{2} r_0^2 \sum_{\mathrm{edge}} (\text{edge length}) \right) \left(\frac{r_{\mathrm{pc}}}{r_0} \right)^2 \\
&= \phi_0 \left(\frac{r_{\mathrm{pc}}}{r_0} \right)^2
\end{aligned}
$$

and thus $r_{\mathrm{pc}} = r_0 \sqrt{\dfrac{\phi_{\mathrm{noise}}}{\phi_0}}$.

References

[1] F. Lipnizki, Beer clarification, in: E. Drioli, L. Giorno (Eds.), Encyclopedia of Membranes, Springer, Berlin, Heidelberg, 2015, pp. 1–2, http://dx.doi.org/10.1007/978-3-642-40872-4_46-1.

[2] C.J. Huang, B.M. Yang, K.S. Chen, C.C. Chang, C.M. Kao, Application of membrane technology on semiconductor wastewater reclamation: A pilot-scale study, Desalination 278 (1) (2011) 203–210, http://dx.doi.org/10.1016/j.desal.2011.05.032.

[3] G. Liu, M. Xiao, X. Zhang, C. Gal, X. Chen, L. Liu, S. Pan, J. Wu, L. Tang, D. Clements-Croome, A review of air filtration technologies for sustainable and healthy building ventilation, Sustain. Cities Soc. 32 (2017) 375–396, http://dx.doi.org/10.1016/j.scs.2017.04.011.

[4] N. Dizge, D.Y. Koseoglu-Imer, A. Karagunduz, B. Keskinler, Effects of cationic polyelectrolyte on filterability and fouling reduction of submerged membrane bioreactor (MBR), J. Membr. Sci. 377 (1) (2011) 175–181, http://dx.doi.org/10.1016/j.memsci.2011.04.048.

[5] C.L. Chang, M.S. Chang, Preparation of multi-layer silicone/PVDF composite membranes for pervaporation of ethanol aqueous solutions, J. Membr. Sci. 238 (1) (2004) 117–122, http://dx.doi.org/10.1016/j.memsci.2004.03.026.

[6] A. Onur, A. Ng, W. Batchelor, G. Garnier, Multi-layer filters: Adsorption and filtration mechanisms for improved separation, Front. Chem. 6 (2018) 417, http://dx.doi.org/10.3389/fchem.2018.00417.

[7] A. Patanaik, R.D. Anandjiwala, L. Boguslavsky, Development of high efficiency particulate absorbing filter materials, J. Appl. Polym. Sci. 114 (1) (2009) 275–280, http://dx.doi.org/10.1002/app.30470.

[8] Y.C. Ahn, S.K. Park, G.T. Kim, Y.J. Hwang, C.G. Lee, H.S. Shin, J.K. Lee, Development of high efficiency nanofilters made of nanofibers, Curr. Appl. Phys. 6 (6) (2006) 1030–1035, http://dx.doi.org/10.1016/j.cap.2005.07.013, Nano Korea 2004 Symposium on NT Challenge.

[9] R.S. Barhate, S. Ramakrishna, Nanofibrous filtering media: Filtration problems and solutions from tiny materials, J. Membr. Sci. 296 (1) (2007) 1–8, http://dx.doi.org/10.1016/j.memsci.2007.03.038.

[10] D. Fong, L.J. Cummings, S.J. Chapman, P. Sanaei, On the performance of multilayered membrane filters, J. Eng. Math. 127 (23) (2021) 1–25, http://dx.doi.org/10.1007/s10665-021-10118-2.

[11] P. Sanaei, L.J. Cummings, Membrane filtration with complex branching pore morphology, Phys. Rev. Fluids 3 (2018) 094305, http://dx.doi.org/10.1103/PhysRevFluids.3.094305.

[12] Y.S. Polyakov, Depth filtration approach to the theory of standard blocking: Prediction of membrane permeation rate and selectivity, J. Membr. Sci. 322 (1) (2008) 81–90, http://dx.doi.org/10.1016/j.memsci.2008.05.012.

[13] J. Jimenez-Martinez, C.F.A. Negre, Eigenvector centrality for geometric and topological characterization of porous media, Phys. Rev. E 96 (2017) 013310, http://dx.doi.org/10.1103/PhysRevE.96.013310.

[14] W. Sun, T. Chen, C. Chen, J. Li, A study on membrane morphology by digital image processing, J. Membr. Sci. 305 (1) (2007) 93–102, http://dx.doi.org/10.1016/j.memsci.2007.07.040.

[15] M.J. Blunt, B. Bijeljic, H. Dong, O. Gharbi, S. Iglauer, P. Mostaghimi, A. Paluszny, C. Pentland, Pore-scale imaging and modelling, Adv. Water Resour. 51 (2013) 197–216, http://dx.doi.org/10.1016/j.advwatres.2012.03.003, 35th Year Anniversary Issue.

[16] G. Sundaramoorthi, M. Hadwiger, M.B. Romdhane, A.R. Behzad, R. Madhavan, S.P. Nunes, 3D membrane imaging and porosity visualization, Ind. Eng. Chem. Res. 55 (12) (2016) 3689–3695, http://dx.doi.org/10.1021/acs.iecr.6b00387.

[17] B. Bijeljic, M.G. Andrew, H. Dong, A.Q. Raeini, M.J. Blunt, P. Mostaghimi, Z. Alhashmi, iRock Technologies, Micro-CT images and networks, 2007, URL https://www.imperial.ac.uk/earth-science/research/research-groups/pore-scale-modelling/micro-ct-images-and-networks/.

[18] A.U. Krupp, I.M. Griffiths, C.P. Please, Stochastic modelling of membrane filtration, Proc. R Soc. A 473 (2200) (2017) 20160948, http://dx.doi.org/10.1098/rspa.2016.0948.

[19] I.M. Griffiths, I. Mitevski, I. Vujkovac, M.R. Illingworth, P.S. Stewart, The role of tortuosity in filtration efficiency: A general network model for filtration, J. Membr. Sci. 598 (2020) 117664, http://dx.doi.org/10.1016/j.memsci.2019.117664.

[20] B. Gu, L. Kondic, L.J. Cummings, A graphical representation of membrane filtration, SIAM J. Appl. Math. 82 (2022) 3, http://dx.doi.org/10.1137/21M1424743.

[21] P.K.P. Reis, M.S. Carvalho, Pore-scale analysis of condensate blockage mitigation by wettability alteration, Energy J. 13 (18) (2020) 4673, http://dx.doi.org/10.3390/en13184673.

[22] X. Tan, D. Rodrigue, A review on porous polymeric membrane preparation. Part I: Production techniques with polysulfone and poly (vinylidene fluoride), Polymers (ISSN: 2073-4360) 11 (7) (2019) 1160, http://dx.doi.org/10.3390/polym11071160.

[23] H. Tao, Spongelike structure, in: D. Enrico, G. Lidietta (Eds.), Encyclopedia of Membranes, Springer Berlin Heidelberg, Berlin, Heidelberg, ISBN: 978-3-662-44324-8, 2016, pp. 1814–1815, http://dx.doi.org/10.1007/978-3-662-44324-8_1954.

[24] J. Yan, K. Dong, Y. Zhang, X. Wang, A.A. Aboalhassanand J Yu, B. Ding, Multifunctional flexible membranes from sponge-like porous carbon nanofibers with high conductivity, Nature Commun. (ISSN: 2041-1723) 10 (1) (2019) 5584, http://dx.doi.org/10.1038/s41467-019-13430-9.

[25] S. Mochizuki, A.L. Zydney, Theoretical analysis of pore size distribution effects on membrane transport, J. Membr. Sci. 82 (3) (1993) 211–227, http://dx.doi.org/10.1016/0376-7388(93)85186-Z.

[26] S. Saksena, A.L. Zydney, Pore size distribution effects on electrokinetic phenomena in semipermeable membranes, J. Membr. Sci. 105 (3) (1995) 203–215, http://dx.doi.org/10.1016/0376-7388(95)00060-P.

[27] D.M. Kanani, W.H. Fissell, S. Roy, A. Dubnisheva, A. Fleischman, A.L. Zydney, Permeability–selectivity analysis for ultrafiltration: Effect of pore geometry, J. Membr. Sci. 349 (1) (2010) 405–410, http://dx.doi.org/10.1016/j.memsci.2009.12.003.

[28] Y. Liu, J. Gao, Y. Ge, S. Yu, M. Liu, C. Gao, A combined interfacial polymerization and in-situ sol-gel strategy to construct composite nanofiltration membrane with improved pore size distribution and anti-protein-fouling property, J. Membr. Sci. 623 (2021) 119097, http://dx.doi.org/10.1016/j.memsci.2021.119097.

[29] N. Taylor, W.J. Ma, A. Kristopeit, S. Wang, A.L. Zydney, Retention characteristics of sterile filters – effect of pore size and structure, J. Membr. Sci. 635 (2021) 119436, http://dx.doi.org/10.1016/j.memsci.2021.119436.

[30] A.L. Zydney, P. Aimar, M. Meireles, J.M. Pimbley, G. Belfort, Use of the log-normal probability density function to analyze membrane pore size distributions: functional forms and discrepancies, J. Membr. Sci. 91 (3) (1994) 293–298, http://dx.doi.org/10.1016/0376-7388(94)80090-1.

[31] K. Xu, W. Wei, Y. Chen, H. Tian, S. Xu, J. Cai, A pore network approach to study throat size effect on the permeability of reconstructed porous media, Water (ISSN: 2073-4441) 14 (1) (2022) 77, http://dx.doi.org/10.3390/w14010077.

[32] Subhadeep Roy, Santanu Sinha, Alex Hansen, Role of pore-size distribution on effective rheology of two-phase flow in porous media, Front. Water (ISSN: 2624-9375) 3 (2021) http://dx.doi.org/10.3389/frwa.2021.709833, URL https://www.frontiersin.org/article/10.3389/frwa.2021.709833.

[33] Y.S. Lin, A.J. Burggraaf, Experimental studies on pore size change of porous ceramic membranes after modification, J. Membr. Sci. (ISSN: 0376-7388) 79 (1) (1993) 65–82, http://dx.doi.org/10.1016/0376-7388(93)85018-R.

[34] Sung-Wook Choi, Yu Zhang, Younan Xia, Three-dimensional scaffolds for tissue engineering: the importance of uniformity in pore size and structure, Langmuir 26 (2010) 19001–19006, http://dx.doi.org/10.1021/la104206h.

[35] U.S. Environmental Protection Agency, Technical fact sheet – 1,4-dioxane, 2017, URL https://www.epa.gov/fedfac/technical-fact-sheet-14-dioxane.