

Stability and collapse of holes in liquid layers with vibrations



Abstract

We investigate experimentally and theoretically the stability and collapse of holes in liquid layers on bounded substrates with the addition of surface vibrations.

It is shown that for a liquid layer with a thickness of the order of the capillary length, a stable hole exists when the hole diameter is bigger than a critical value d_c .

Consequently, a further increase of the liquid volume causes the hole to collapse. It is found that d_c increases with the size of the container, but decreases with the magnitude of the surface vibrations [1].

Motivation

Control of the evolution of vibrated wetting fluids on the micro-scale. Example: Smoothly adding liquid to a vibrated film containing a hole until an unstable state is reached [1].

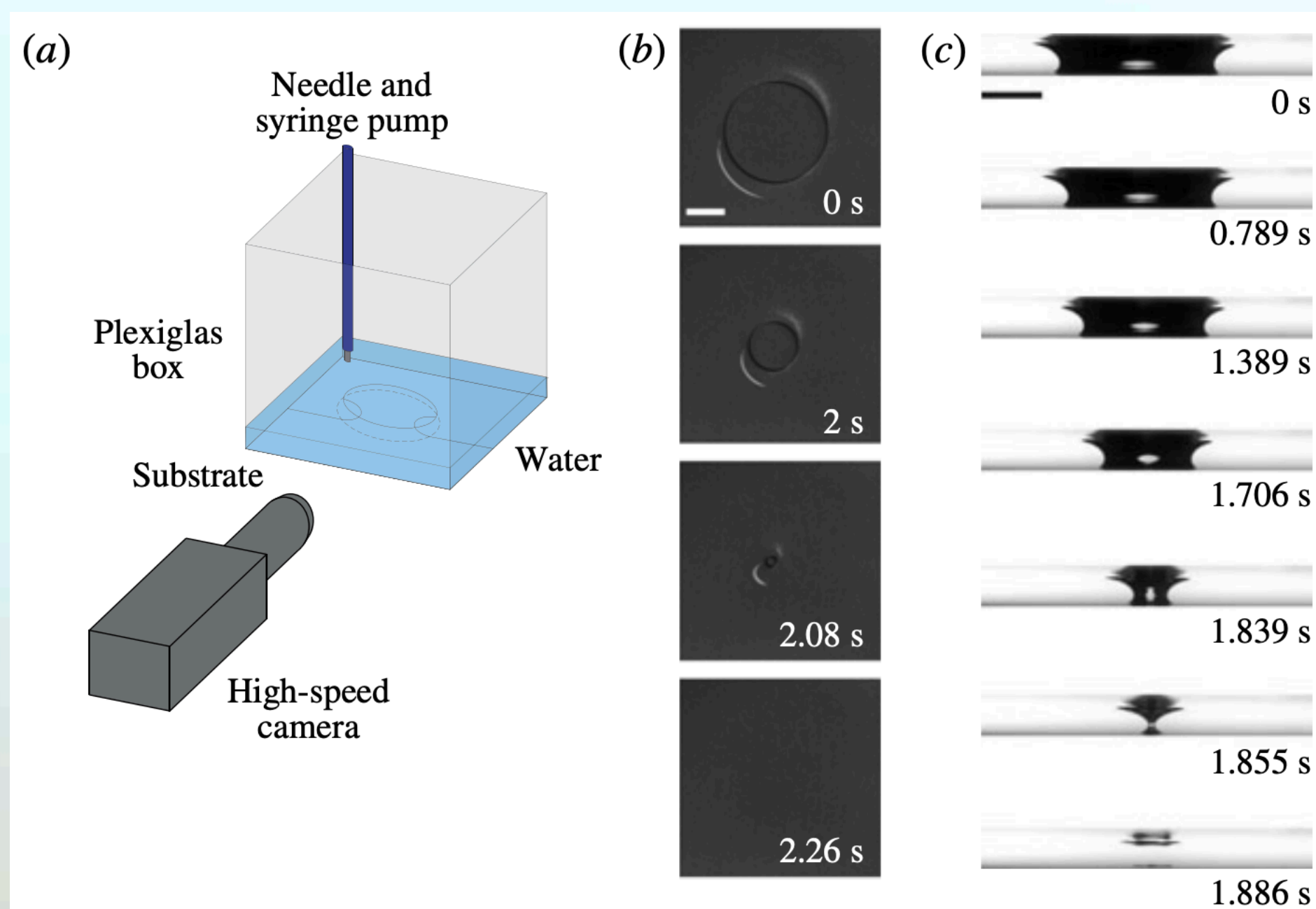


Fig. 1: (a) Schematic illustration of the experimental set-up with a needle connected to a syringe pump controlling the flow rate. Time lapse images of the collapse of a hole at the centre of a liquid layer on a bounded superhydrophobic Al plate are shown in top (b) and side (c) view, which are results from two different experiments.

Main assumptions

- Long wave theory describing thin film dynamics is used [2].
- Direction normal to the plane must be shorter than any in plane dimensions
- The slope of the free surface is small
- We are working in an axis-symmetric coordinate system
- Neumann conditions were used at the left and right boundary along with zero flux condition on the left but fixed non-zero flux on the right

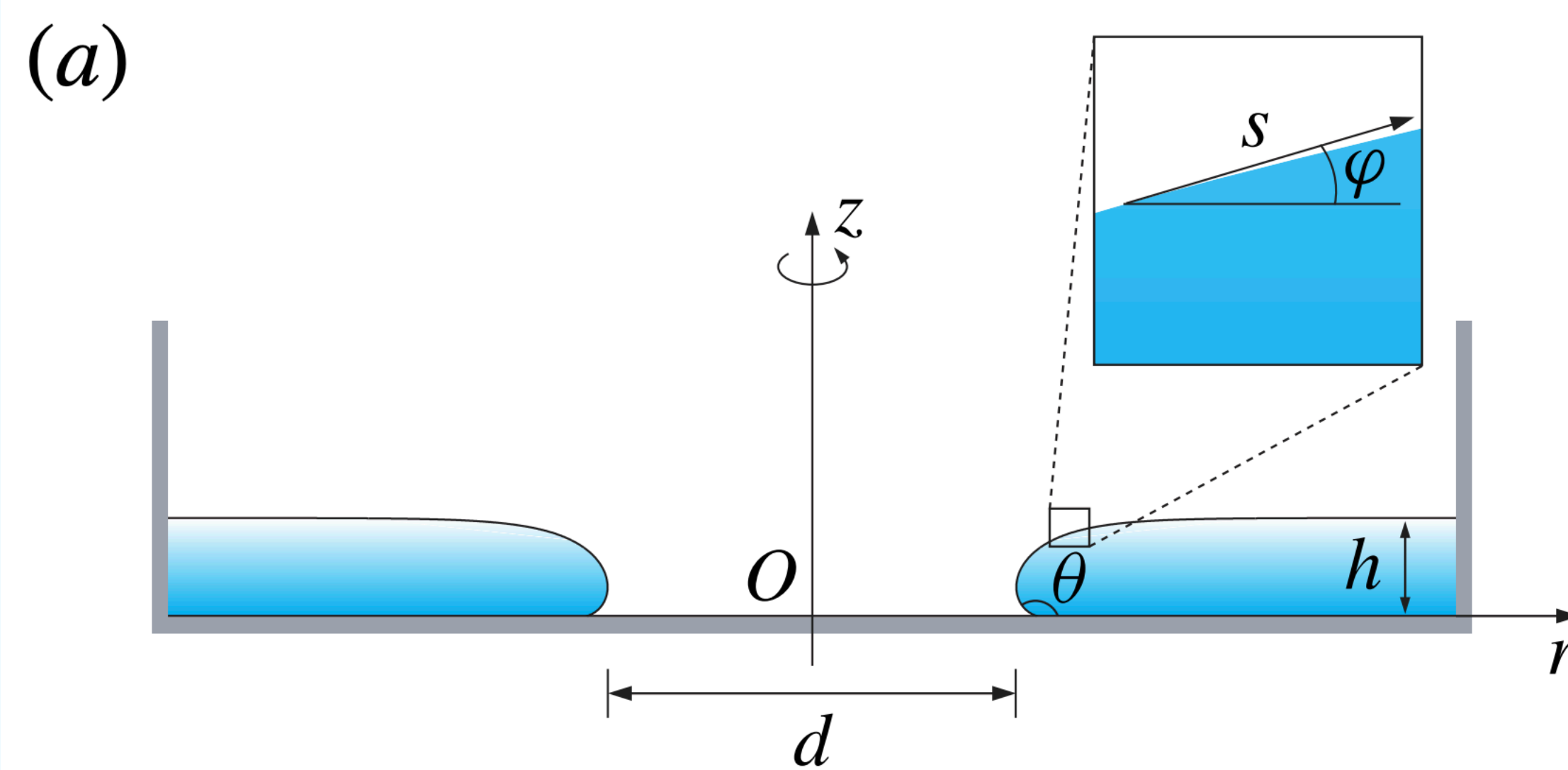


Fig. 2: Sketch of the modeling problem.

Equations

$$\frac{\partial h}{\partial t} = \frac{1}{r} \left[r h^3 \left((1 + \hat{C} \cos(\omega t)) h_r - K f' h_r - T_r \right) \right]_r$$

$$\hat{C} = \frac{\bar{C}}{g} = \frac{A \omega^2}{g} = \frac{4 \pi^2 A f^2}{g}$$

$$K = \frac{\kappa a}{\gamma} = \frac{(1 - \cos(\theta))(m - 1)(n - 1)}{h_*(n - m)}$$

$$f(h) = \left(\frac{h_*}{h} \right)^n - \left(\frac{h_*}{h} \right)^m$$

$$T = \frac{1}{r} (r h_r)_r$$

$$x_c = \left(\frac{\gamma h_c}{\rho g} \right)^{1/3}$$

$$t_c = \frac{3 \mu x_c^4}{\gamma h_c^3}$$

Results

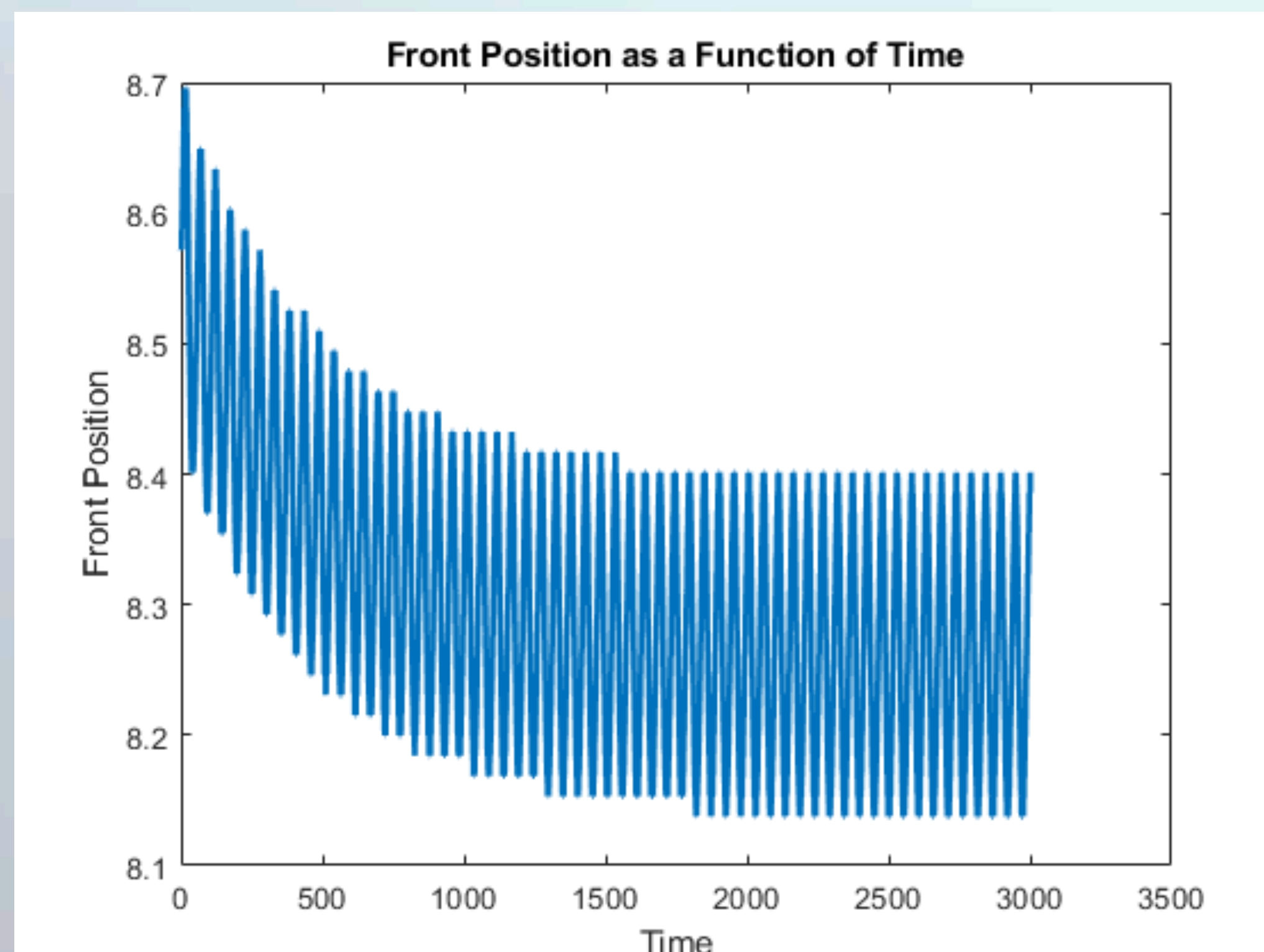


Fig. 3: Fluid front evolution as a function of time for $\hat{C} = -2$

The fluid front oscillates with the same frequency as the surface vibrations and approaches a new stable radius size as time progresses.

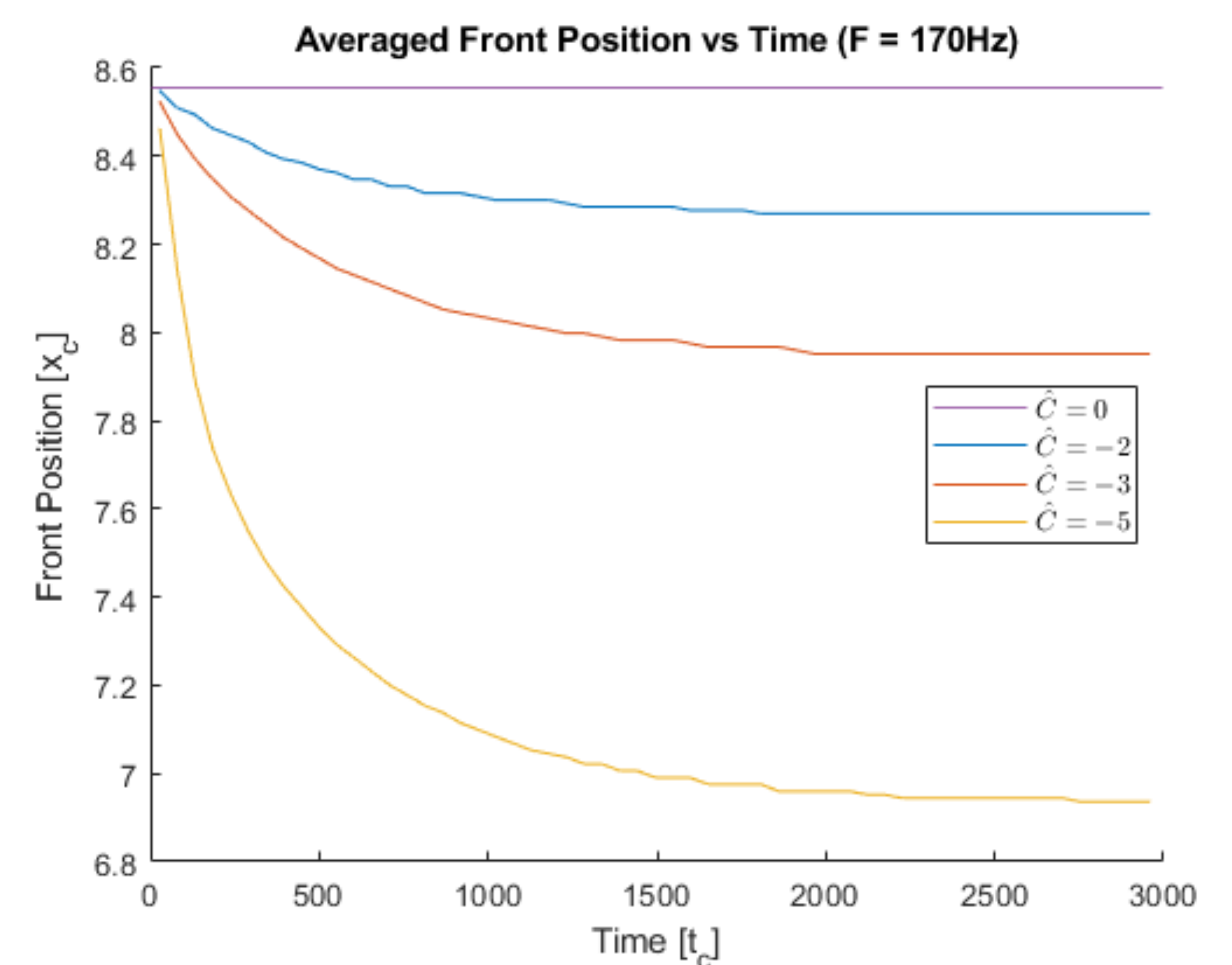


Fig. 4: Averaged front position as a function of time for decreasing \hat{C} values

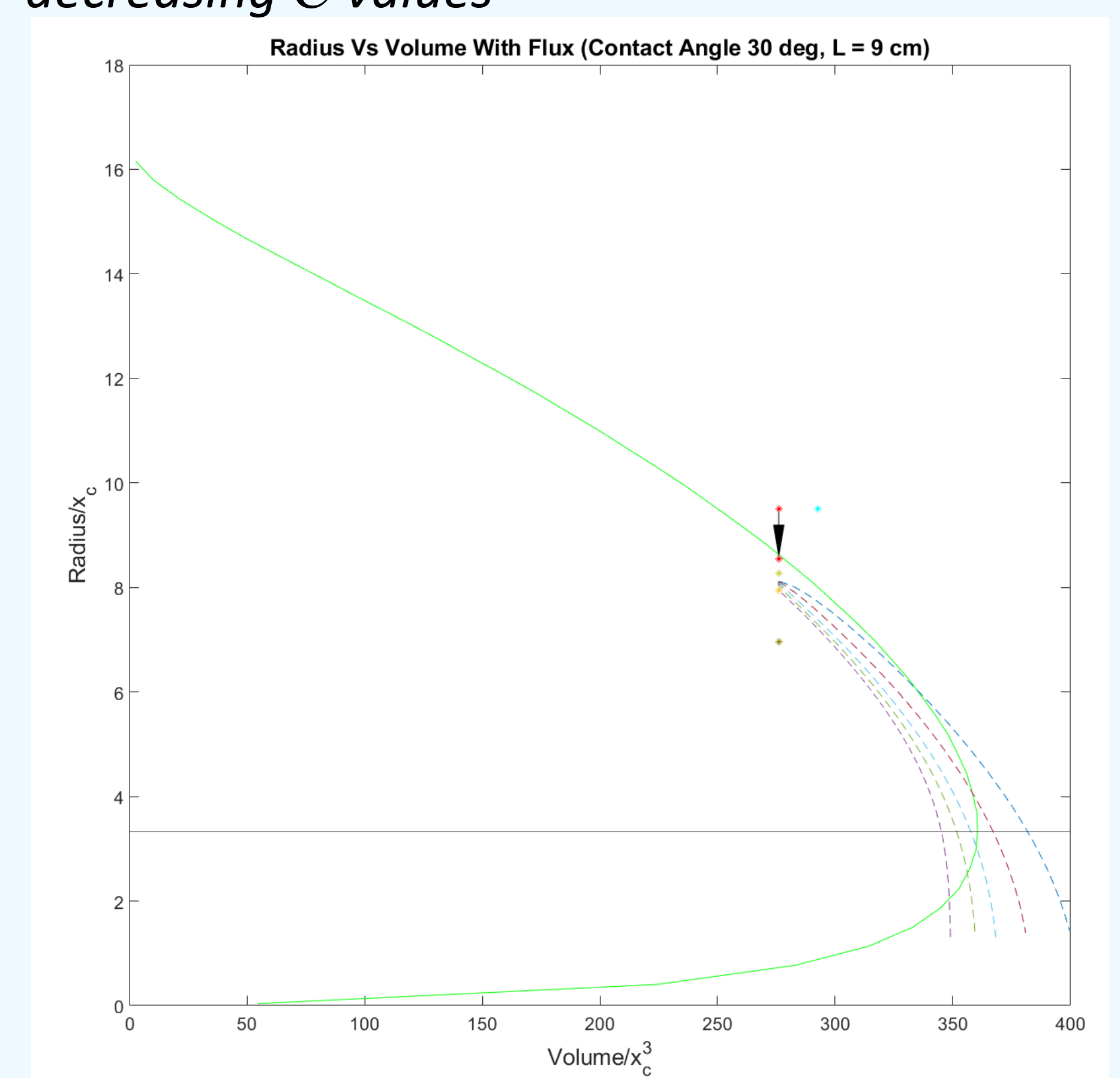


Fig. 5: Fluid front stability plot for $\hat{C} = -3$ and decreasing values of supplied flux

Summary

Vertical surface vibrations can cause a decrease in the stable hole size of holes in liquid layers.

As flux is supplied at decreasing rates we can find the new critical diameter for a given value of \hat{C} .

This new d_c is shown to be less than the critical diameter without vibrations.

Future Work: Experiments involve larger values of Reynolds number so we would like to include inertial terms in our simulations.

References

- [1] Lv, Cunjing, et al. "Stability and Collapse of Holes in Liquid Layers." *Journal of Fluid Mechanics*, vol. 855, 2018, pp. 1130–1155.
- [2] Kondic, L. "Instabilities in Gravity Driven Flow of Thin Fluid Films." *SIAM Review*, vol. 45, no. 1, 2003, pp. 95–115.