

# A Free-Streamline Model for a Liquid Film in a Fast Flow

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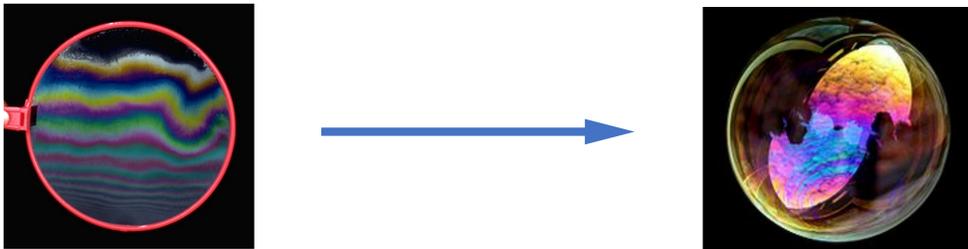
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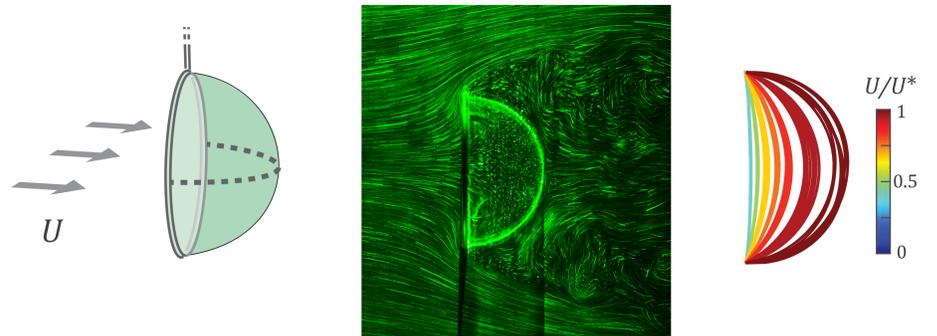
## Motivation

- How does a blown film become a bubble?
- What are the hydrodynamic events that trigger necking down and pinching off?

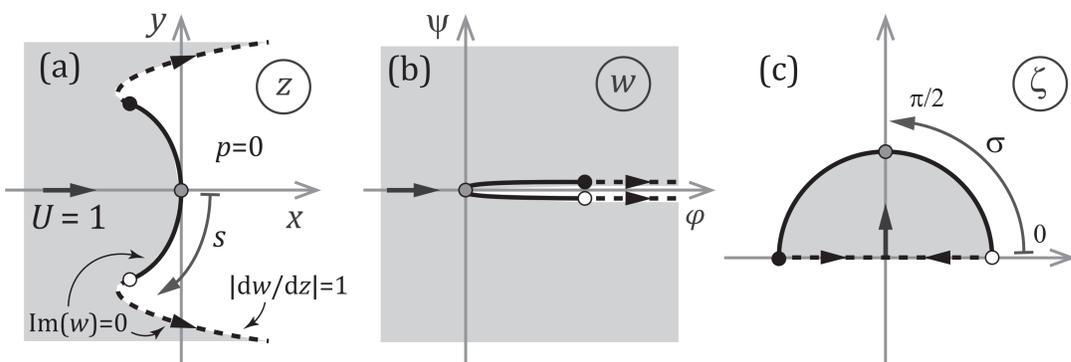


## Experimental setup

- Consider a liquid-liquid system: an oil film in a water tunnel.
- The film inflates as the flow speed  $U$  is increased, and ruptures at a critical speed  $U^*$ .



## Free-streamline theory: a 1D film in a 2D flow



$$\frac{d\theta}{d\sigma} = -K\eta \sinh[\tau(\sigma)] \sin 2\sigma, \quad \theta(\pi/2) = -\pi/2,$$

$$\frac{dz}{d\sigma} = -\frac{K}{2} e^{i\theta} e^{-\tau} \sin 2\sigma, \quad \text{Im}[z(0)] = -1,$$

$$\tau = \mathcal{H}[\theta].$$

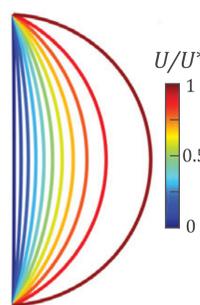
- Flow (gray) is assumed to be inviscid and irrotational.
- Free streamlines (dashed curves) have unknown shape, and shed tangentially from the pinning points.
- The wake (white) is assumed to have uniform pressure.
- Film shape (black curve) is determined by the balance of hydrodynamic pressure and surface tension.
- The potential flow is found using conformal mapping.
- Let  $\Omega = i \log(dw/dz)$ , where  $w(z)$  is the velocity potential.
- Define  $\Omega = \theta(\sigma) + i\tau(\sigma)$  on the film, and  $w(\zeta) = (K/8)(\zeta + \zeta^{-1})^2$ .
- $\Omega \in \mathbb{R}$  for  $\zeta \in \mathbb{R}$  (free streamline). Schwarz reflection principle implies that  $\Omega$  may be extended to the unit disc in the  $\zeta$ -plane:
 
$$\Omega(\zeta) = \sum_k a_k \zeta^k \quad \text{where } a_k \in \mathbb{R}.$$
 Thus,  $\tau = \mathcal{H}[\theta]$ , where  $\mathcal{H}$  is the Hilbert transform.
- Governing equation is obtained by combining
 
$$p = 2\gamma\kappa \quad \text{(Young-Laplace equation)}$$

$$p + \frac{1}{2}\rho|\mathbf{u}|^2 = \frac{1}{2}\rho U^2 \quad \text{(Bernoulli's Law)}$$
 where  $p$  is the hydrodynamic pressure,  $\gamma$  the surface tension,  $\kappa$  the film curvature,  $\rho$  the fluid density, and  $\mathbf{u}$  the flow velocity.
- The Weber number is  $\eta = \rho U^2 R / 4\gamma$ , where  $R$  is the ring radius.
- Equation is solved numerically using a Broyden method, with  $\eta$  as the continuation parameter.

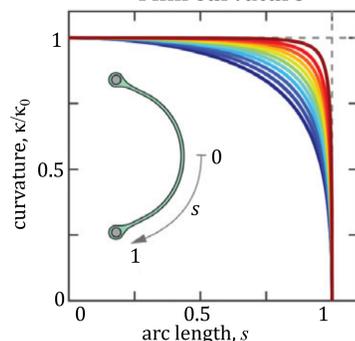
## Results

- Highly inflated films have nearly uniform curvature, except near the pinning points.

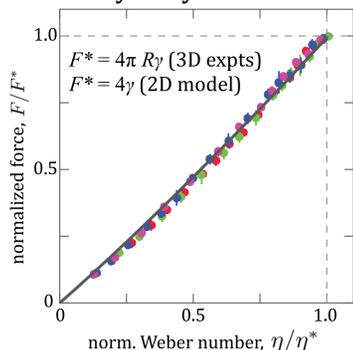
### Predicted film shapes



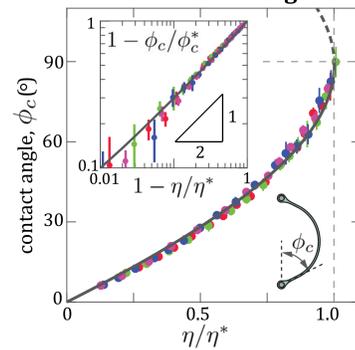
### Film curvature



### Hydrodynamic force

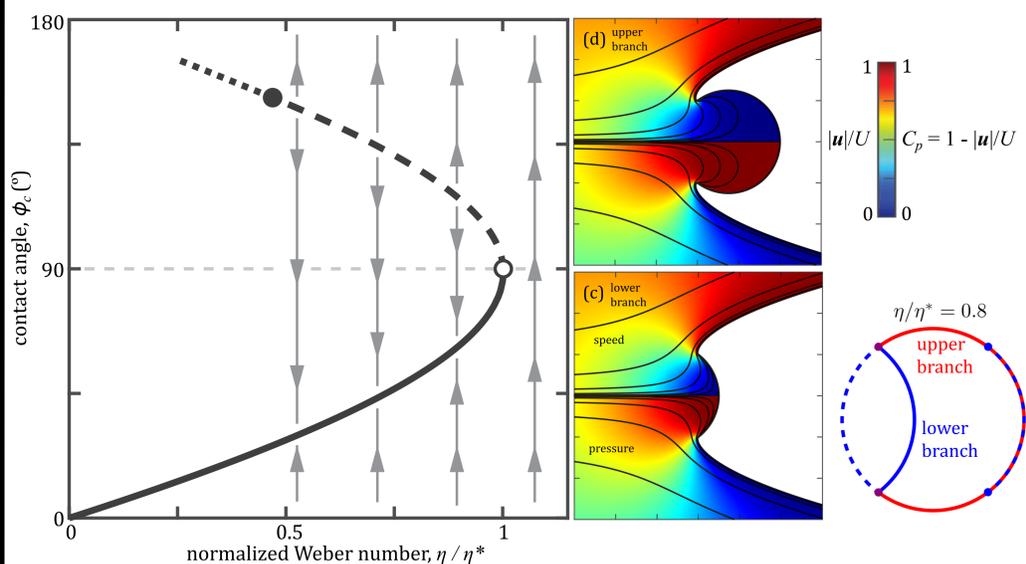


### Film contact angle



- Theory exhibits favorable agreement with experiment.

## Model prediction: over-inflated equilibrium shapes



- **Stable shape:** Inward deformations lower the Laplace pressure due to the decrease in curvature, so the external pressure dominates and restores the shape.
- **Unstable shape:** Inward perturbations increase curvature, leading to further collapse of the film.

## Conclusions

- Rupture results from a maximal surface tension force that can resist imposed fluid pressures.
- Rupture is associated with the critical point of a saddle-node bifurcation.
- Future work: what conditions determine the different fates of rupture vs. birthing of a bubble?

## References & Acknowledgments

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- Reference: Ganedi *et al.*, "Equilibrium states and their stability for liquid films in fast flows," *Physical Review Letters* **121**, 094501 (2018).